

# Sensitivity of $G_0$ and stress-strain relation of geomaterials to grain shape and surface roughness

Masahide Otsubo<sup>1,2#</sup>, Yang Li<sup>2</sup>, and Reiko Kuwano<sup>2</sup>

<sup>1</sup>Public Works Research Institute, 1-6 Minamihara, Tsukuba, Ibaraki 305-8516, Japan

<sup>2</sup>The University of Tokyo, Institute of Industrial Science, 4-6-1 Komaba, Meguro-ku, Tokyo 153-8505, Japan

<sup>#</sup>Corresponding author: [ootsubo-m573cm@pwri.go.jp](mailto:ootsubo-m573cm@pwri.go.jp)

## ABSTRACT

It is empirically known that the packing property and mechanical responses of cohesionless granular materials are influenced by grain shape. For example, the attainable range of void ratio depends on grain shape; angular materials tend to exhibit greater friction angles. Besides, shear wave velocity ( $V_s$ ) and small-strain shear modulus ( $G_0$ ) of sphere assemblies are affected by surface roughness. However, consensus has yet to be reached on the combined effect of grain shape and surface roughness on  $G_0$  and stress-strain relation. This contribution aims to evaluate the shape-roughness combined effect on the strain-dependent mechanical responses of granular materials. Three groups of glass beads having different grain shapes and a silica sand were used, and their grain surfaces were roughened through a systematic procedure using a milling machine. In total, eight materials were subjected to triaxial compression after measurement of  $V_s$ . The experimental results reveal that the stress-strain relation of angular particles is remarkably affected by the surface roughness, whereas the roughness effect on the stress-dependent variation in  $G_0$  is limited for angular particles.

**Keywords:** Grain shape; surface roughness; stiffness; strength.

## 1. Introduction

Granular materials are composed of a large number of particles and their overall mechanical response can be considered as a consequence of complicated interactions of grains. It is empirically known that both packing characteristics and stress-strain relation of cohesionless granular materials are influenced by grain shape (Cho et al. 2006; Altuhafi et al. 2016). It has also been demonstrated that shear wave velocity ( $V_s$ ) and thus small-strain shear modulus ( $G_0$ ) of sphere assemblies tend to be lower for spherical particles having rougher surfaces (Santamarina and Cascante 1998; Sharifipour and Dano 2006; Otsubo et al. 2015). However, consensus has yet to be reached on the potential effect of surface roughness on friction angles. Further, there has been little consideration on the combined effect of grain shape and surface roughness on the  $G_0$  and stress-strain relation.

This contribution aims to evaluate the shape-roughness combined effect on the strain-dependent mechanical responses of granular materials focusing on both stiffness and strength characteristics in experiments.

## 2. Materials

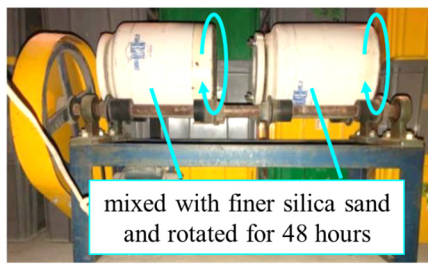
Three groups of glass beads (GB) with different grain shapes and a silica sand (SS), having low uniformity coefficients ( $C_u$ ), were used (Table 1). The three glass beads are spherical (SGB), deformed (DGB) and angular (AGB). Their as-supplied surfaces were roughened using a milling machine (Fig. 1) following Dutta et al. (2020). Each material and a fine silica sand ( $D_{50}=0.15$  mm) were mixed in a ceramic jar having a ratio of 1:2 by mass, and rotated for 48 hours with about 1.5 Hz. The as-supplied natural silica sand (NSS) was also roughened to prepare rough-surface silica sand (RSS). Grain sizes and shapes of the as-supplied materials were measured using a QICPIC apparatus (Altuhafi et al. 2013), and they were assumed unchanged through the gentle milling process. The extreme void ratios ( $e_{min}$  and  $e_{max}$ ) for the tested materials are given in Table 2. The average and root-mean square surface roughness ( $S_a$  and  $S_q$ ) were quantified using Gaussian filter method (Li et al. 2021).

**Table 1.** Particle properties ( $G_s$ : specific gravity) and shape parameters ( $S$ : sphericity;  $C_x$ : convexity;  $AR$ : aspect ratio, Altuhafi et al. 2013)

ID	$G_s$	$D_{50}$ [mm]	$C_u$	$S$	$C_x$	$AR$
<b>SSGB</b>	2.50	0.496	1.15	0.940	0.974	0.975
<b>SDGB</b>	2.50	1.040	1.17	0.844	0.982	0.784
<b>SAGB</b>	2.50	0.746	1.31	0.786	0.959	0.631
<b>NSS</b>	2.64	0.525	1.46	0.857	0.936	0.720

**Table 2.** Surface roughness and extreme void ratios

ID	$S_a$ [nm]	$S_q$ [nm]	$e_{min}$	$e_{max}$
<b>SSGB</b>	38.9	67.6	0.589	0.694
<b>RSGB</b>	114.3	173.5	0.606	0.696
<b>SDGB</b>	97.5	154.8	0.494	0.666
<b>RDGB</b>	171.7	232.0	0.509	0.683
<b>SAGB</b>	75.0	127.2	0.734	1.176
<b>RAGB</b>	148.0	224.7	0.745	1.170
<b>NSS</b>	424.3	680.0	0.719	1.122
<b>RSS</b>	479.9	795.4	0.716	1.060



**Figure 1.** Milling machine used to roughen grain surface.

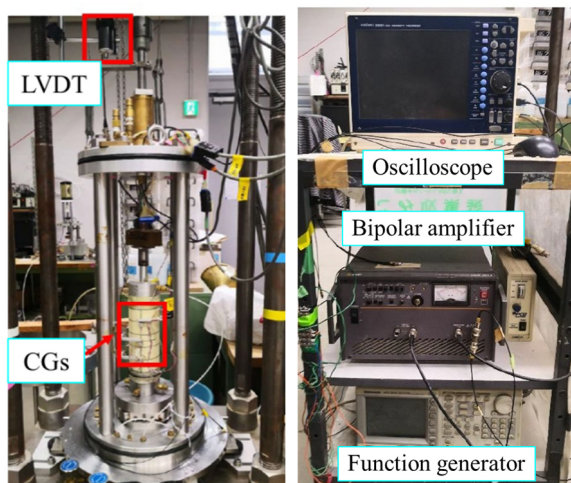
### 3. Test method

The adopted test method in this contribution is described in detail in Li (2022) that broadly follows the approach in Dutta et al. (2020; 2021).

#### 3.1. Sample preparation

Cylindrical samples having 75 mm in diameter and 150 mm in height ( $H$ ) were prepared using a split mould. Dry materials were carefully poured in air using a funnel and gentle side tapping was given to the mould to densify the sample to vary initial void ratio ( $e$ ) and relative density ( $D_r$ ) in a systematic manner. Each material was prepared at least two distinct density levels.

Samples were then subjected to an initial isotropic stress ( $p'_0$ ) of 50 kPa under dry and drained conditions by applying negative pressure, where  $V_s$  was measured as explained in the following section. For the stress levels above 50 kPa, cell pressure was applied to the sample under dry and drained conditions. One axial displacement sensor (linear variable differential transducer, LVDT) and three radial displacement transducers (clip gauges, CGs) were placed to monitor the deformation of the sample under dry and drained conditions (Fig. 2).



**Figure 2.** Test equipment used: (left) triaxial apparatus, (right) shear wave measurement system.

#### 3.2. Shear wave measurement

At the initial isotropic state of 50 kPa, shear wave measurement was conducted using shear plates (Fig. 3). The system comprises of shear (S-) and compression (P-) type piezoelectric plates, having 20 mm in diameter and

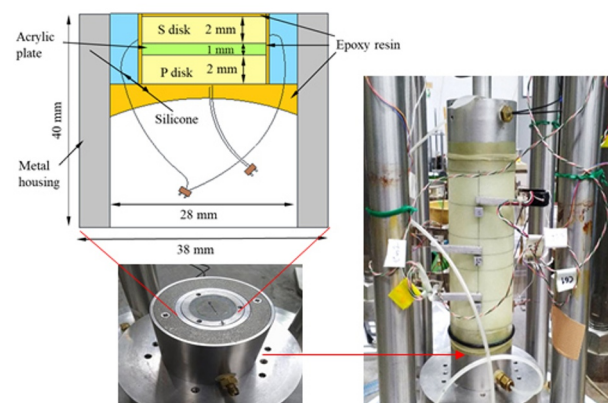
2 mm in thickness, inside a hollow stainless-steel cylinder, which is supported by epoxy glue and silicone (Fig. 3) as presented in Dutta et al. (2021). This contribution utilizes the S-wave data (from shear-type plate) to determine shear wave velocity ( $V_s$ ) measured in the vertical direction.

Two conventional methods for estimating the travel time from time-domain responses are peak-to-peak ( $ptp$ ) and start-to-start ( $sts$ ) methods. Yamashita et al. (2009) recommend that input frequency ( $f_{in}$ ) should be selected such that  $ptp$  and  $sts$  methods give comparable  $V_s$ . Following Dutta et al. (2020; 2021), this contribution uses a single-period sinusoidal form having  $f_{in}$  of 7 kHz as an inserted wave. The travel time ( $T_T$ ) of shear waves was determined using the  $ptp$  method. Typical wave signals measured for the tested materials are presented in Li (2022). To amplify the received signals, the inserted voltage was amplified to  $\pm 70$  V via a bipolar amplifier (Fig. 2).  $V_s$  was calculated as  $H/T_T$ .

Assuming an isotropic and homogeneous condition, the following relation can be used to estimate  $G_0$  from  $V_s$  measured in the vertical direction.

$$G_0 = \rho_d V_s^2 \quad (1)$$

where  $\rho_d$  is the dry density of the sample and is expressed as  $G_s \rho_w / (1 + e)$ ;  $\rho_w$  is the density of water.



**Figure 3.** Schematic of shear plates (disk-shaped transducers) in metal housing inside top cap and pedestal (Dutta et al. 2021).

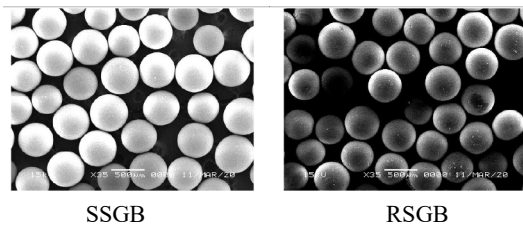
#### 3.3. Triaxial compression

After conducting shear wave measurement at  $p'_0 = 50$  kPa, the sample was monotonically sheared by increasing the axial strain up to  $\epsilon_a$  of about 15% (defined as the residual state in this study). The minor principal stress ( $\sigma'_3$ ) was kept at 50 kPa while increasing the axial stress ( $\sigma'_1$ ). A constant strain rate of 0.0006 %/s was adopted to record shear wave signals at target axial strains. During the shearing, deviator stress,  $q = \sigma'_1 - \sigma'_3$ , and mean effective stress,  $p' = (\sigma'_1 + 2\sigma'_3)/3$  were recorded. In the following section, the evolution of deviator stress ratio ( $q/p'$ ) against  $\epsilon_a$  is discussed. Besides, mobilized peak and residual friction angles ( $\phi_{peak}$  and  $\phi_{res}$ ) for each material are provided. While the data are not presented here, Li (2022) reports the evolutions of volumetric strain and S- and P-wave velocities during the shearing process.

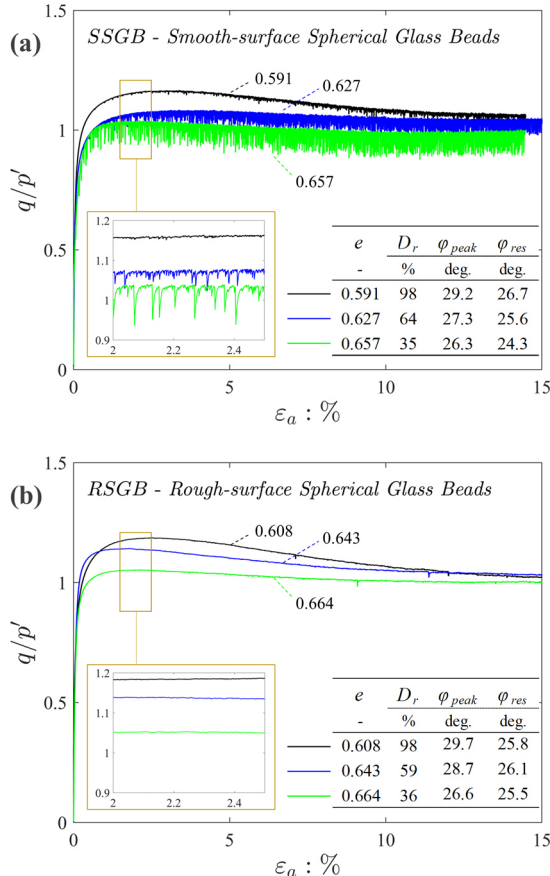
## 4. Stress-strain relation

### 4.1. Spherical glass beads (SGB)

For smooth-surface and rough-surface spherical glass beads (SSGB and RSGB, Fig. 4), stress-strain curves for three samples having different  $e$  are plotted in Fig. 5. Regarding the peak strength, RSGB samples having a rough surface exhibit slightly greater  $\varphi_{peak}$  values under the equivalent  $D_r$ , where  $e$  is slightly higher for RSGB. Dutta et al. (2020) reported that under an equivalent  $e$ , rougher materials exhibit larger  $\varphi_{peak}$  than the smooth equivalent; this agrees with the result in Fig. 5. More dilation is evident for rougher materials (Li 2022). The  $\varphi_{res}$  values are similar between SSGB and RSGB; this trend agrees with the finding in Cavarretta et al. (2010). SSGB with higher  $e$  exhibit fluctuation in  $q/p'$ ; when this is significant,  $\varphi$  values were calculated considering the median value of the fluctuation amplitude.



**Figure 4.** SEM images of spherical glass beads having smooth surface (SSGB) and roughened surface (RSGB).



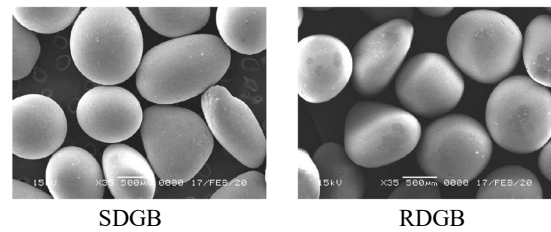
**Figure 5.** Stress-strain curves for (a) smooth-surface spherical glass beads (SSGB) and (b) rough-surface spherical glass beads (RSGB) at  $\sigma'_3 = 50$  kPa.

### 4.2. Deformed glass beads (DGB)

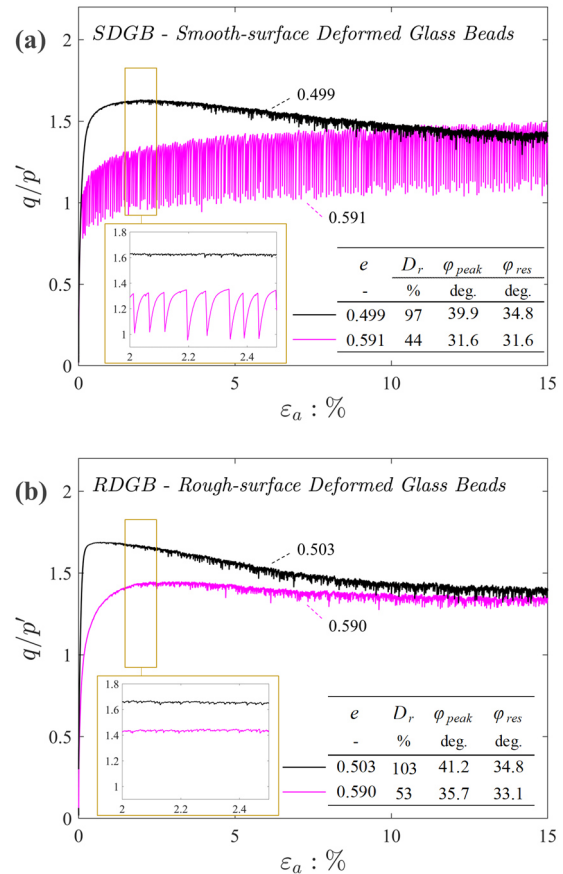
Smooth-surface deformed glass beads (SDGB, Fig. 6) were produced in the same process as SSGB, where particles having slightly deformed shapes were collected as SDGB in the manufacturing company.

Fig. 7 shows that  $q/p'$  for SDGB is clearly greater than that for SSGB or RSGB, indicating a greater contribution of grain shape than surface roughness. The effect of surface roughness appears in a similar manner with the comparison between SSGB and RSGB (Fig. 5). That is, the  $\varphi_{peak}$  values of RDGB are slightly greater than SDGB, whereas the  $\varphi_{res}$  values are similar irrespective of the magnitude of surface roughness.

For the loose SDGB sample (Fig. 7a), fluctuation in  $q/p'$  is more significant than SSGB (Fig. 5a), probably due to the larger grain sizes in SDGB (Table 1). Nevertheless, the fluctuation is removed in RDGB, probably due to the prevention of inter-particle slip.



**Figure 6.** SEM images of deformed glass beads having smooth surface (SDGB) and roughened surface (RDGB).



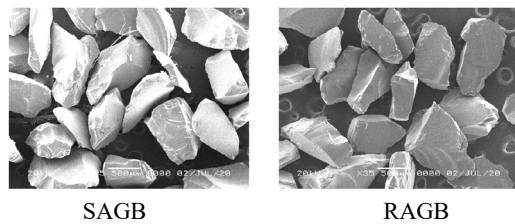
**Figure 7.** Stress-strain curves for (a) smooth-surface deformed glass beads (SDGB) and (b) rough-surface deformed glass beads (RDGB) at  $\sigma'_3 = 50$  kPa.

### 4.3. Angular glass beads (AGB)

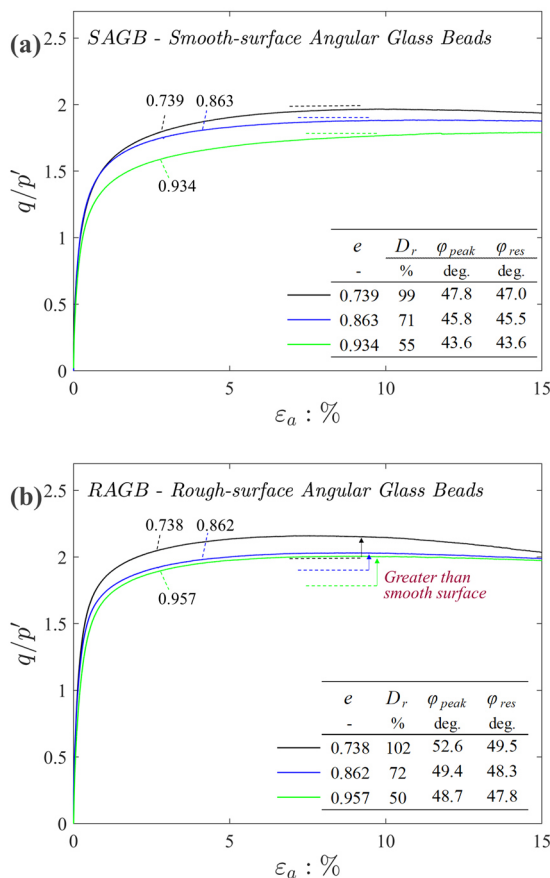
Smooth-surface angular glass beads (SAGB, Fig. 8) were produced by crushing spherical glass beads in the manufacturing company. Thus, SAGB and SSGB have the same physical properties except for their grain shapes.

Fig. 9 illustrates that the  $\varphi_{peak}$  and  $\varphi_{res}$  values of AGB samples are much greater than the SGB and DGB materials despite higher  $e$  in the AGB samples. This highlights a remarkable role of grain shape. However, the initial slope of the  $q/p' - \varepsilon_a$  relation of AGB samples appears lower than that for SGB and DGB samples, probably due to the difference in  $e$  or grain shape.

Contrary to SGB and DGB samples, the effect of surface roughness on AGB samples is obvious. For a similar  $e$  or  $D_r$ , both the  $\varphi_{peak}$  and  $\varphi_{res}$  values are larger for RAGB, suggesting that the effect of surface roughness is more remarkable for angular materials.



**Figure 8.** SEM images of angular glass beads having smooth surface (SAGB) and roughened surface (RAGB).

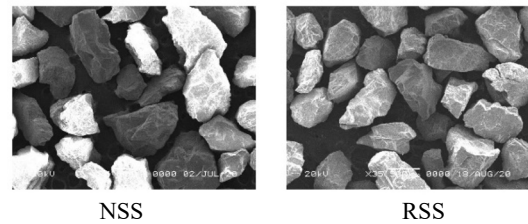


**Figure 9.** Stress-strain curves for (a) smooth-surface angular glass beads (SAGB) and (b) rough-surface angular glass beads (RAGB) at  $\sigma_3' = 50$  kPa. Dashed lines indicate the difference between SAGB and RAGB samples at equivalent densities.

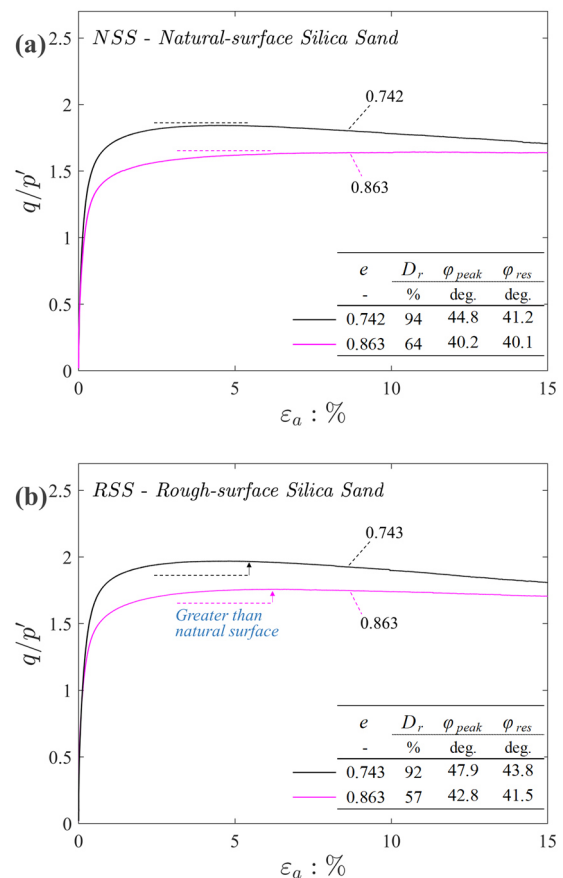
### 4.4. Silica sand (SS)

The natural-surface silica sand selected has been widely used in soil mechanics research in Japan. The main component is  $SiO_2$ , similar with quartz sand. In the manufacturing process, the silica sand was produced by crushing a larger mass. This process is similar with the SAGB material, and thus their shapes are indeed similar (Figs. 8 and 10). It is noteworthy that the  $e_{min}$  and  $e_{max}$  also similar between the two materials (Table 2). The natural-surface was made rougher in the same milling process adopted for RSSB, RDGB and RAGB.

The overall stress-strain curves of NSS and RSS are similar with SAGB and RAGB, respectively, where the  $\varphi_{peak}$  of SS samples are slightly lower than AGB samples. Under the equivalent  $e$ , the effect of surface roughness is evident; RSS exhibits greater  $\varphi_{peak}$  and  $\varphi_{res}$  values than NSS (Fig. 11).



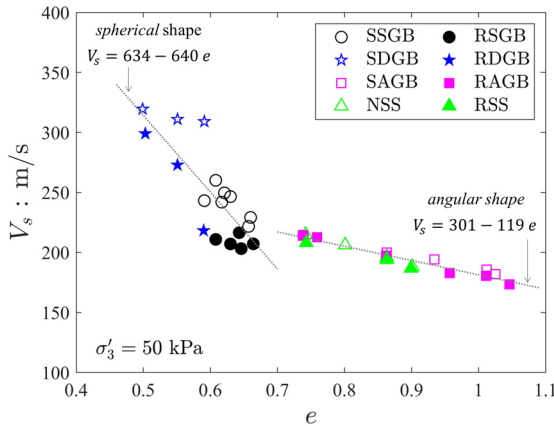
**Figure 10.** SEM images of silica sands having natural (as-provided) surface (NSS) and roughened surface (RSS).



**Figure 11.** Stress-strain curves for (a) natural-surface (as-provided) silica sand (NSS) and (b) rough-surface silica sand (RSS) at  $\sigma_3' = 50$  kPa. Dashed lines highlight the difference between NSS and RSS samples at equivalent densities.

## 5. Small-strain shear modulus

At  $p'_0 = 50$  kPa,  $V_s$  was measured at various  $e$  for the eight materials (Fig. 12). Referring to Table 2, each material has different  $e_{min}$  and  $e_{max}$ ; SGB and DGB cannot attain  $e > 0.7$ , while AGB and SS cannot achieve  $e < 0.7$ . Fig. 12 shows fitting lines based on regression analyses for two groups separately, namely *spherical shape* (SGB and DGB) and *angular shape* (AGB and SS).



**Figure 12.** Relationship between  $V_s$  and  $e$  at  $\sigma'_3 = 50$  kPa (Fitting lines are plotted for *spherical shape* (SGB and DGB) and *angular shape* (AGB and SS), separately).

Referring to Hardin and Richart (1963), the  $V_s - e$  relation can be expressed linearly as:

$$V_s = A_1 - A_2 e \quad (2)$$

$$B = A_1/A_2 \quad (3)$$

Following the approach in Hardin and Richart (1963), the  $B$  value in Eq. (3) can be used in a void ratio correction function for  $G_0$  as:

$$f(e) = (B - e)^2 / (1 + e) \quad (4)$$

The  $B$  value for SSGB alone is 1.26, similar to 1.28 for borosilicate beads (Otsubo et al. 2018) and 1.44 for the same glass bead material having a larger size (Dutta et al. 2020). However, the dataset is scattered in Fig. 12, especially for the DGB samples. Therefore, two separate fitting lines including SGB and DGB (i.e. *spherical shape*) and including AGB and SS (i.e. *angular shape*) were adopted for the later discussion, giving  $B = 0.991$  and 2.52, respectively. The  $B$  value for the group of *angular shape* is similar to 2.97 proposed for angular grains (dry crushed quartz sands) in Hardin and Richart (1963). The overall trend in Fig. 12 also captures the transition observed between round grains and angular grains in Hardin and Richart (1963).

From Eq. (1),  $G_0$  can be calculated as a function of  $e$  and  $V_s$ . For each material, one sample was tested at  $p'_0 = 25, 50, 100, 200$  and 400 kPa to explore the stress dependent variation of  $G_0$  (Fig. 13). To remove the effect of void ratio on  $G_0$  at different stress levels, Eq. (4) can be used. To express the stress-dependent evolution of  $G_0$ , the following expression is often used.

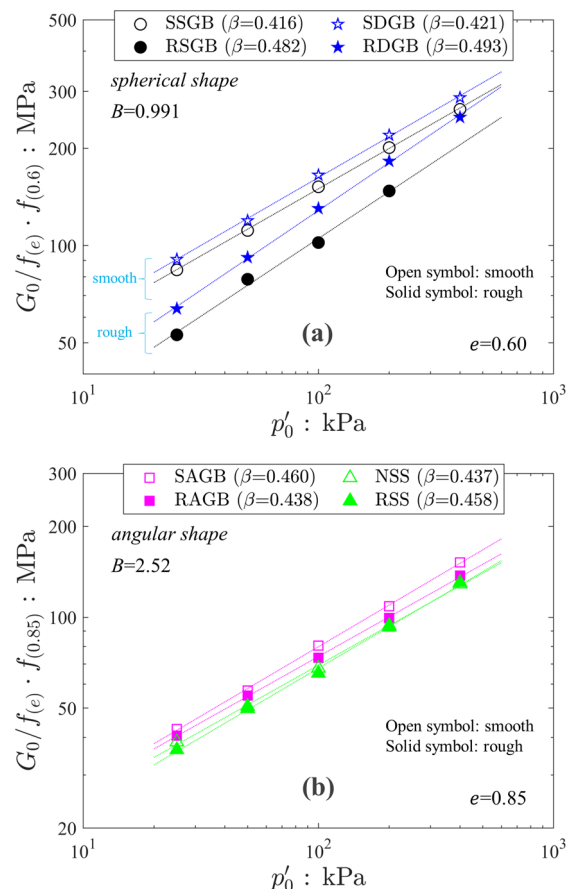
$$G_0 = \alpha f(e) p_0'^\beta \quad (5)$$

where  $\alpha$  is the material constant and  $\beta$  is the stress exponent. Fig. 13 compares the evolution of  $G_0$  for a reference void ratio of 0.6 and 0.85, respectively, for *spherical shape* and *angular shape* groups.

The magnitude of  $\alpha$  depends on  $f(e)$ , and thus  $\alpha$  depends on  $B$ . Referring to Fig. 12 and Eq. (3), this contribution uses two different  $B$  values for *spherical shape* (SGB and DGB) and *angular shape* (AGB and SS); thus, comparison on  $\alpha$  alone is not very meaningful.

On the other hand,  $\beta$  has been reported sensitive to the surface roughness for spherical particles (e.g. Yimsiri and Soga 2000). Fig. 13a shows that  $\beta$  increases from 0.42 (smooth) to 0.49 (rough) for SGB and DGB samples. Otsubo and O'Sullivan (2018) reported experimentally that  $S_q/D_{50}$  can be related to  $\beta$ , ranging from 0.38 to 0.54 for spherical glass beads. In this study, SSGB has a non-negligible  $S_q/D_{50}$  than Otsubo and O'Sullivan (2018); thus, giving a slightly larger  $\beta$  value of 0.416.

Referring to Fig. 13a, DGB samples exhibit similar trends with SGB samples, whereas AGB and SS samples give different trends in which  $\beta$  is not sensitive to the surface roughness ( $\beta = 0.44$  to 0.46) (Fig. 13b). RAGB gives even lower  $\beta$  than SAGB, contrary to the observation for SGB and DGB. McDowell and Bolton (2001) reported that  $\beta$  is 0.5 when a conical-shape asperity is in contact with a flat plate; sharp edges of angular grains could cause such a conical-shape contact. It can be concluded that the effect of surface roughness on  $G_0$  is limited for angular materials.



**Figure 13.** Evolution of  $G_0$  with increasing  $p'_0$  (a) *spherical shape* (SGB and DGB) for  $e=0.6$  (b) *angular shape* (AGB and SS) for  $e=0.85$  ( $\beta$  is obtained from Eq. (3)).

## 6. Conclusions

This contribution explored the influence of grain shape, surface roughness and their combination on the mechanical response of cohesionless granular materials. Based on a series of triaxial compression tests along with shear wave measurements at the initial isotropic stress state, the following conclusions can be drawn:

- Small-strain shear modulus ( $G_0$ ) of spherical materials is sensitively affected by the surface roughness where  $G_0$  tends to be lower for rougher surface material; however, the stress exponent ( $\beta$ ) is higher for materials with rougher surface (Fig. 13a).
- However, the effect of surface roughness is limited for angular materials, i.e.  $G_0$  of angular materials is not sensitive to the difference in surface roughness.  $\beta$  is as high as 0.460 even for the smooth-surface angular glass beads (SAGB) (Fig. 13b).
- The range between  $e_{max}$  and  $e_{min}$  is different between spherical materials and angular materials. This makes it difficult to discuss the effect of grain shape on  $G_0$  systematically. For the attainable range of  $e$ , however, spherical particles (with lower  $e$ ) tend to give greater  $V_s$  and thus  $G_0$  (Fig. 12).
- Peak and residual strengths (i.e.  $\varphi_{peak}$  and  $\varphi_{res}$ ) are largely influenced by grain shape, much significant than the effect of surface roughness.
- Increased surface roughness leads to an increase in the  $\varphi_{peak}$  and  $\varphi_{res}$  values. This is more remarkable for angular materials than spherical particles (Figs. 9 and 11). Particularly, this effect is limited for the comparison between SSGB and RSDB (Fig. 5).
- As a side effect, significant fluctuation observed in the stress-strain curves for the smooth-surface spherical particles (SSGB and SDGB) disappear when their surfaces were roughened. This indicates that inter-particle slip can be prevented due to increased surface roughness.

The present contribution has been developed based on Li (2022), focusing on uniformly graded granular materials tested at a relatively low stress level ( $\sigma'_3 = 50$  kPa). Therefore, further research is needed to draw a clearer picture of the combined effect of grain shape and surface roughness on the strain-dependent mechanical response of geomaterials.

## Acknowledgements

This work was supported by JSPS KAKENHI Grant Number 22K14322, and JSPS Bilateral Joint Research Projects (JPJSBP120195701) with The Royal Society (IEC\R3\183026).

## References

Altuhafi, F., O'Sullivan, C., Cavarretta, I. 2013. "Analysis of an image-based method to quantify the size and shape of

- sand particles." *J Geotech Geoenviron.* 139, no. 8, 1290–1307.  
[https://doi.org/10.1061/\(ASCE\)GT.1943-5606.0000855](https://doi.org/10.1061/(ASCE)GT.1943-5606.0000855)
- Altuhafi, F., Coop, M., Georgiannou, V. 2016. "Effect of particle shape on the mechanical behavior of natural sands." *J Geotech Geoenviron.* 142, no. 12, 04016071.  
[https://doi.org/10.1061/\(ASCE\)GT.1943-5606.0001569](https://doi.org/10.1061/(ASCE)GT.1943-5606.0001569)
- Cavarretta, I., Coop, M., O'Sullivan, C. 2010. "The influence of particle characteristics on the behaviour of coarse grained soils." *Géotechnique* 60, no. 6, 413–423.  
<https://doi.org/10.1680/geot.2010.60.6.413>
- Cho, G. C., Dodds, J., Santamarina, J. C. 2006. "Particle shape effects on packing density, stiffness, and strength: natural and crushed sands." *J Geotech Geoenviron.* 132, no. 5, 591–602.  
[https://doi.org/10.1061/\(ASCE\)1090-0241\(2006\)132:5\(591\)](https://doi.org/10.1061/(ASCE)1090-0241(2006)132:5(591))
- Dutta, T., Otsubo, M., Kuwano, R., O'Sullivan, C. 2020. "Evolution of shear wave velocity during triaxial compression." *Soils Found.* 60, no. 6, 1357–1370.  
<https://doi.org/10.1016/j.sandf.2020.07.008>
- Dutta, T., Otsubo, M., Kuwano, R. 2021. "Effect of shearing history on stress wave velocities of sands observed in triaxial compression tests." *Soils Found.* 61, no. 2, 541–548.  
<https://doi.org/10.1016/j.sandf.2021.02.002>
- Hardin, B. O., Richart Jr., F. E. 1963. "Elastic wave velocities in granular soils." *J Soil Mech Found Divis.* 89, no. 1, 33–66. <https://doi.org/10.1061/JSFEAQ.0000493>
- McDowell, G., Bolton, M., 2001. Micro mechanics of elastic soil. *Soils Found.* 41, no. 6, 147–152.  
[https://doi.org/10.3208/sandf.41.6\\_147](https://doi.org/10.3208/sandf.41.6_147)
- Santamarina, J. C., Cascante, G. 1998. "Effect of surface roughness on wave propagation parameters." *Géotechnique* 48, no. 1, 129–136.  
<https://doi.org/10.1680/geot.1998.48.1.129>
- Sharifipour, M., Dano, C. 2006. "Effect of grains roughness on waves velocities in granular packings." In *Proceedings of the First Euro Mediterranean Symposium in Advances on Geomaterials and Structures*, edited by F. Darve, 123–128. Hammamet, Tunisia. <https://hal.science/hal-01008097>
- Otsubo, M., O'Sullivan, C., Sim, W. W., Ibraim, E. 2015. "Quantitative assessment of the influence of surface roughness on soil stiffness." *Géotechnique* 65, no. 8, 694–700.  
<https://doi.org/10.1680/geot.14.T.028>
- Otsubo, M., O'Sullivan, C. 2018. "Experimental and DEM assessment of the stress-dependency of surface roughness effects on shear modulus". *Soils Found.* 58, no 3, 602–614.  
<https://doi.org/10.1016/j.sandf.2018.02.020>
- Li, Y. "Effects of particle morphology on mechanical responses of granular materials", *PhD thesis*, The university of Tokyo, 2022.
- Li, Y., Otsubo, M., Kuwano, R., Nadimi, S. 2021. "Quantitative evaluation of surface roughness for granular materials using Gaussian filter method". *Powder Technology* 388, 251–260.  
<https://doi.org/10.1016/j.powtec.2021.04.068>
- Yamashita, S., Kawaguchi, T., Nakata, Y., Mikami, T., Fujiwara, T., Shibuya, S. 2009. "Interpretation of international parallel test on the measurement of  $G_{max}$  using bender elements." *Soils Found.* 49, no. 4, 631–650.  
<https://doi.org/10.3208/sandf.49.631>
- Yimsiri, S., Soga, K. 2000. "Micromechanics-based stress-strain behaviour of soils at small strains." *Géotechnique* 50, no. 5, 559–571. <https://doi.org/10.1680/geot.2000.50.5.559>