

# Study of the motion of two-phase systems based on fluid mechanics and physical modelling of fluids

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**Abstract.** The movement of bipolar symptoms is discovered through nature and practice. Analyses have shown that the lack of models has limited one to laboratory experiments with limited empirical images of their results. In the last decades, a basic stage has been reached with a large amount of laboratory work and experimental studies, where data and analyses have been systematically processed. In recent years, many works have included reviews of the current state of the problem and the development of individual processes, mainly using empirical methods. In this work, an analytical study of hydrodynamic processes in two-phase flow is presented and subsequently applied under field conditions. This approach avoids the limitations of the field and can be easily used to determine parameters of interest in engineering problems. Our solution is based on the results of mental theory.

**Keywords:** Shear stress, Dense trace gas content, Two-phase flow, Gas-liquid mixture.

## 1. INTRODUCTION

A large number of flows have shown that two-phase flow follows all the basic laws of hydrodynamics; it should be noted that the heterogeneous fluid equations in their systems are more complex and numerous than in single-phase flow in aquatic systems. Methods for analysing one-dimensional patterns are divided into several categories depending on the complexity of the amount of information necessary to describe the system for two-phase flow.

It is well known that a variety of equations and methods can be used to describe real physical processes. As the most common and developed method of describing physical processes in the motion of two-phase networks in fluid mechanics, the continuity assumptions, laws and methods of mechanics of continuous media are used. According to this fluid mechanics, gas-liquid mixtures can be considered as a special branch of the mechanics of continuous media. It is worth noting that a two-phase liquid system can be considered as a medium because the physical properties of the fields are different; the curtains can change in both spaces. The preservation of Zelkoni-Zekonov created during this time and in mechanics is These laws are conservation of mass, momentum impulse, angular momentum (momentum impulse Eperia) and entropy balance.

## 2. PROBLEM STATEMENT

To solve multinuclear problems, models are needed. As the model description indicates, this means that the use of models to study physical processes allows the corresponding processes to be studied with the help of modelling in order to obtain reliable information or to test hypotheses. The base is based on the model of A. A. Armand.

This paper, entitled "Motion study of a two-phase system based on hydrodynamic and hydrophysical simulations", deals with oil from electrochemical methods, bathing water with other methods and problems, the impact of combined methods, homogeneous and stratified heterogeneity, and field-based formation body processes.

## 3. THE MAIN PART

Firstly, calculate the liquid viscosity:

$$\varepsilon_{\mu} = \varepsilon_{\mu}(1 - \varphi)^{-1.5}$$

Or

$$\varepsilon_{\mu} = \mu_l \frac{d\mu_l}{d\tau} (1 - \varphi)^{-1.5} \quad (1)$$

$\mu_l$ -Viscosity;

$\varepsilon_{\mu}$ -tangential stress;

$\varepsilon_1$ -additional stresses in the fluid in motion

Therefore, for the sake of equilibrium, Vasrigodorindra will be calculated during normal motion, and the sum of

all these pre-existing external forces, according to the laws of mechanics, will give the results

$$P_1 - P_2 - T = 0$$

Or

$$P_1 \pi \tau^2 - P_3 \pi \tau^2 + \mu \frac{d\mu_l}{d\tau} 2N_2 L (1 - \varphi)^{-1.5} = 0 \quad (2)$$

Subtracting  $N_2^2$  gives:

$$P_1 - P_2 = \Delta P = -\mu_l \frac{d\mu_l}{d\tau} 2L (1 - \varphi)^{-1.5} \quad (3)$$

It's available:

$$\Delta P \frac{\tau(1 - \varphi)^{1.5}}{2L} d\tau = -\mu_l d\mu_l$$

Or

$$\Delta P \frac{(1 - \varphi)^{1.5}}{2L\mu_l} \tau d\tau = -d\mu_l \quad (4)$$

$P_1$  and  $P_2$ -represent the densities of the gas and liquid phases;

$\frac{d\mu_l}{d\tau}$ -Gradients of homogenised liquids;

$\varphi$ -component;

Length

This paper investigates the parabolic phenomenon elucidated by the equations. The constant integrals of boundary conditions, determined by the following velocity on the pipe wall, i.e., the velocity of a partially rotating bullet, i.e.,  $\tau = \tau_0 \mu_l = 0$ . The integral of this equation provides the velocity distribution across the cross-section of uniform gas-liquid flow in a circular pipe:

$$0 = -\frac{\tau_0^2}{\varphi \mu_l} \frac{\Delta \varphi (1 - \varphi)^{1.5}}{L} + C$$

Or

$$0 = \frac{\tau_0^2}{\varphi \mu_l} \frac{\Delta P}{L} (1 - \varphi)^{1.5} \quad (5)$$

As a result:

$$\mu_l = -\frac{\tau^2}{\varphi \mu_l} \frac{\Delta \varphi (1 - \varphi)^{1.5}}{L} + \frac{\tau_0^2}{\varphi \mu_l} \frac{\Delta \varphi (1 - \varphi)^{1.5}}{L}$$

Or

$$\mu_l = \frac{\Delta P}{\varphi \mu_l L} (1 - \varphi)^{1.5} (\tau_0^2 - \tau^2)$$

So eventually get:

$$\mu_l = \frac{\Delta P \tau_0^2}{\varphi \mu_l L} (1 - \varphi)^{1.5} (1 - \frac{\tau^2}{\tau_0^2}) \quad (6)$$

This equation determines the rate at which the fluid changes in a uniform flow. If  $\varphi = 0$  then we have

$$\mu_l = \frac{\Delta P \tau_0^2}{\varphi \mu_l L} (1 - \frac{\tau^2}{\tau_0^2}) \quad (7)$$

The equation, initially derived by Stokes in 1867, expresses the ground distribution of uniform fluid velocity across the cross-section of a pipe under linear motion conditions. Hence, it is referred to as the Parabolic Stokes Law.

he maximum and minimum speeds, determined at:

$$\mu_{maxl} = \frac{\Delta P (1 - \varphi)^{1.5}}{\varphi \mu_l L} \tau_0^2 \quad (8)$$

T, respectively, establish the ratio between the velocity at the maximum speed point and the maximum speed on the wheel axis, denoted as  $\mu_l$ -dimensional velocity, using

$$\frac{\mu_l}{\mu_{maxl}} = 1 - \frac{\tau^2}{\tau_0^2} \quad (9)$$

From the equation, it is evident that under laminar conditions across all active layers of a circular tube, the distribution of non-target velocities remains uniform. Consequently, we can assert that in all scenarios of liquid and gas laminar motion within gas-liquid mixture systems, regardless of fluid viscosity and thus Reynolds number, all influences on fluid motion are similar when non-dimensional velocity values are equal. When this phenomenon bears resemblance to Reynolds' individual significance, it is deemed self-similar. In such instances, the non-followable actual composition becomes its own model, enabling the study of two-phase system motion under conditions where actual velocities decrease manifold. This characteristic of similar phenomena in theoretical models developed by your nation is self-similar.

Using the velocity formula for any point in the aforementioned uniform heat flow cross-section, the formula for determining the flow velocity of liquid family fluids can be ascertained. Considering the content discussed in the flow cross-section, we select a ring whose axis coincides with that of the pipeline, having internal alignment in the radial direction. The area of this ring is  $dF = 2\pi\tau d\tau$ . Taking into account the aforementioned scenario, we use the area to determine the basic flow through this ring:

$$dQ = \mu 2\pi\tau d\tau \quad (10)$$

Upon inserting velocity values into the slots of the table, we obtain:

$$\begin{aligned} dQ &= \mu 2\pi\tau d\tau \\ &= \frac{\Delta \varphi \tau_0^2}{\varphi \mu_l L} (1 - \varphi)^{1.5} (1 - \frac{\tau^2}{\tau_0^2}) 2\pi\tau d\tau \\ &= \frac{\Delta P \tau_0^2 \pi}{2\mu_l L} (1 - \varphi)^{1.5} (\tau d\tau - \frac{\tau^3}{\tau_0^2} d\tau) \\ &= \frac{\pi \Delta \varphi \Delta P}{2\mu_l L} (1 - \varphi)^{1.5} \frac{\tau_0^2}{4} \tau_0^2 \end{aligned}$$

Or

$$Q = \frac{\pi \Delta P}{8\mu_l L} (1 - \varphi)^{1.5} \quad (11)$$

If we assume  $\varphi = 0$ , then we have:

$$Q = \frac{\pi \Delta P}{8\mu_l L} \tau_0^4 \quad (12)$$

Ammonium phosphate. This formula is a consumption formula derived from experiments conducted by our adversary Poiseuille. He studied the motion of capillaries in 1840, applying it to the motion problems of circulatory systems.

Triggered by the Usakh Hydraulic Guide, it is well known that large  $Q = \delta \pi \tau_0^2$ -pae- determines the speed of pulling the DNA armor box through the centerline.

From your speed and the equation, you can determine the average

$$\delta = \frac{Q}{\pi \tau_0^2} = \frac{\pi \Delta P}{8\mu_l L} \tau_0^4 (1 - \varphi)^{1.5} \frac{1}{\pi \tau_0^2} \quad (13)$$

Based on in-situ oxidation of strata, homogeneous and layered heterogeneity, impact of combined methods, other methods and issues, active water with electrochemical utilization, Luental Micros, Edelir. Diakmish, with an error of 0.2%, this pressure gauge is shown to be valid.

Afterwards, we obtain:

$$\delta = \frac{1}{8\mu_1} \frac{\Delta P}{L} \tau_0^2 (1 - \varphi)^{1.5} \quad (14)$$

From the equation for the average velocity of an animal body, we can derive the hydraulic contraction law, which determines the pressure loss during lateral movement of the liquid phase.

$$8\mu_1 \tau_1 L = \Delta P \tau_0^2 (1 - \varphi)^{1.5} \quad (15)$$

Using the diameter instead of the radius, we have:

$$8\mu_1 \tau_1 L = \Delta P \frac{D^2}{L} (1 - \varphi)^{1.5} \quad (16)$$

Or

$$\Delta P = \frac{8 \cdot 4\mu_1 \tau_1 L}{D^2 (1 - \varphi)^{1.5}} \quad (17)$$

We take the buyer's liquid density given on the left and the wallet part as

$$h_l = \frac{\Delta P}{r} = \frac{32\mu_1 \tau_1 L}{r D^2 (1 - \varphi)^{1.5}} \quad (18)$$

This equation is the mathematical expression of the hydraulic resistance law. If  $\varphi = 0$ , then we have a formula for homogeneous fluid.

Thus, the product of the numerator and denominator on the right is given. For equation (18), we find:

$$h_l = \frac{32\mu_1 \delta_l L}{r D^2 (1 - \varphi)^{1.5}} \frac{2\tau_1}{2\tau_1} = \frac{64\mu_1 L \delta_l^2}{r \tau D D^2} \frac{1}{2 (1 - \varphi)^{1.5}} \quad (19)$$

By changing the maximum density  $r = \delta g$  of the kinematic  $\mu_l = \nu_l \delta$  flatness specific all-

$$h_l = \frac{64\nu_l L \tau_1^2}{\delta_l D D^2 2g} \frac{1}{(1 - \varphi)^{1.5}} \quad (20)$$

In the long-range equation, we introduce the concept of the Reynolds number:

$$h_l = \frac{64 L \tau_1^2}{Re D^2 2g} \frac{1}{(1 - \varphi)^{1.5}} \quad (21)$$

The true skin displays a loss of dependence on the Reynolds scaffold.

It noted the relevance of uniform flow of gas-liquid coatings within pipes of equal cross-section in this section. Losses are inversely proportional to the given averages, and depend on their diameter and pipe diameter, multiplied by non-plinny gargoder jelly on both sides, not meaning, purification ( $\delta g$ ). this will Pozher Zeri pressure.

$$\Delta p_l = \frac{64 L \tau_1^2}{Re D^2 2g} \delta \frac{1}{(1 - \varphi)^{1.5}} \quad (22)$$

Formula  $\frac{64}{Re}$  corresponds to a coefficient representing resistance, and according to the discrete petyak  $\lambda = \frac{64}{Re}$  division,  $\mu$  is called the coefficient premya, so the loss can be divided by  $\Delta P = \lambda \frac{L \tau_1^2}{D^2} \delta \frac{1}{(1 - \varphi)^{1.5}}$ .

If there are  $\varphi = 0$ , then the Daren-Weisbach formula for homogeneous liquids in the hydraulic nator loss with respect to the ratio of speed is determined as

$$\xi = \frac{h_l}{\frac{\tau_1^2}{2g} \frac{1}{(1 - \varphi)^{1.5}}} = \lambda \frac{L}{D} \quad (23)$$

It is called Paramébur. This jurisprudential route is built on resistance.

$$\xi = \frac{h_l 2g (1 - \varphi)^{1.5}}{\tau^2} = \frac{\Delta P 2g}{g \delta} = 2E\mu \quad (24)$$

The fluid dynamics alterations in Butiko liquid-liquid mixtures introduce complexity to Turbov Atto. In such scenarios, considering the hydraulic resistance coefficient in the Darcy-Weisbach formula is sufficient for single-phase liquids. The perimeter of the pipeline stands as a crucial parameter for energy loss or flow in both liquid-liquid and gas-liquid mixtures. Therefore, this coefficient becomes indispensable for resolving engineering issues related to fluid motion.

Extensive theoretical and laboratory analyses reveal that due to the intricate nature of turbulent motion, average velocity must substitute instantaneous velocity in engineering calculations. Thus far, a theoretical relationship between the two perimeters remains elusive. Considering this, the theoretical solution to this problem remains incomplete.

Many scientists and researchers are actively exploring this issue. It's not feasible to determine the exchange of resistance between the service and venture capital purely based on theory, as the primary influence is determined by the roughness of the pipeline. The factors influencing the roughness of the pipe wall are diverse and variable, not constant, and can vary under different conditions. Roughness is a major factor causing variations in the resistance coefficient, which has a more realistic impact on actual flow. This parameter contributes to the formation of vortices near the pipe wall and is one of the additional reasons for energy loss in hydraulic resonance and dynamic flow. It has been found by laboratory studies that the roughness of pipe walls creates vortices near their surfaces and appears to be responsible for additional vortex drag and energy losses, e.g., in flows in homogeneous liquids, water, and two-phase mixtures, such as gas-liquid mixtures. It has been experimentally demonstrated that the drag coefficient for purging during turbulent motion depends not only on the Reynolds number but also on the relative roughness. Since the main cache in our model is dominated by the conditions of liquid flow, the consideration of its physical properties in the walk model is the main part of the calculations.

The analyses show that, basically, the Blasius formula is most applicable for homogeneous fluid motion in solid pipes, while for rough pipes the A.D. Altitul formula is more applicable. ( $\lambda = \frac{64}{Re}$ )

According to the theory of Blasius de Pey and the view of A. D. Altitul:

$$\lambda = 0.11 \left( \frac{64}{Re} + \frac{L}{D} \right)^{0.25}$$

Based on this model, the error between laboratory experiments and calculated values is small. Therefore, it is reasonable to recommend the technique for practical applications.

## 4. Conclusions

1) A model has been developed to investigate the fundamental properties of gas properties through laboratory experiments.

2) In the exploration of this model, such an approach was proposed - by analysing the panorama of pressure losses in horizontal pipelines, where the main liquid is oil.

3) This method allows to determine the main indicators of pipeline drive for transporting two-phase systems, i.e., liquid (oil) and gas (air).

In the discussion, a model on the behaviour of liquid fluids in horizontal pipes was presented. As it is well known, physical modelling is a common approach in scientific research. Often, for phenomena of interest, scientists identify similar phenomena by conducting experiments in the laboratory with small or large models, which usually have the same physical properties as the actual object. Physical models are used to study the properties of actual objects.

In this work, it is proposed to study two-dimensional fluid sweeping surfaces in horizontal pipes while considering the general laws of hydrodynamics for homogeneous fluids.

Liquids and gases in pipelines play a special role in the actual transport process. This makes possible the expression of the same protection principles based on which the corresponding protection measures can be established.

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